

Peculiarities of Radiative Loss Profiles and Radiation-Condensation Instability in the T-11M Tokamak Discharges with a Lithium Limiter

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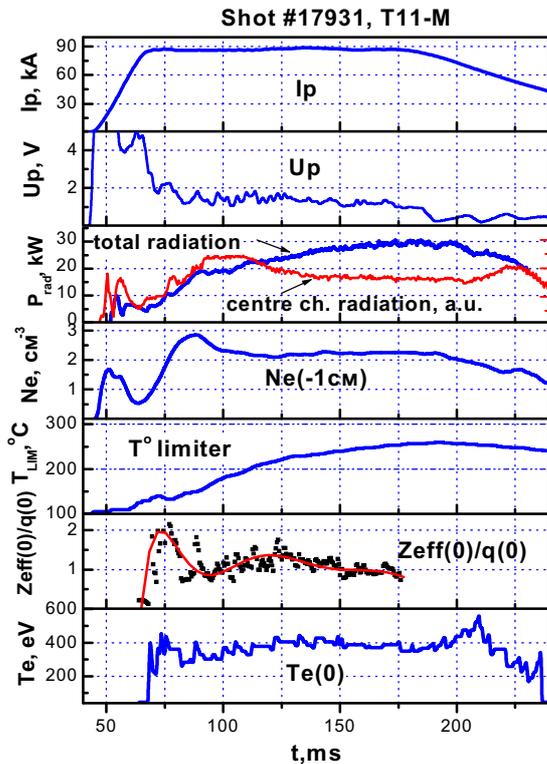


Fig.1. T-11M traces with a thin Li limiter. In this work the reason of this phenomena are considered which, as we have clarified, are related to the radial allocation of radiative losses in the plasma core.

In the last experiments on a tokamak T-11M with thin lithium limiter [1] the quasi-stationary thermal modes of the limiter operation were obtained (Fig.1). At rather low (90...100°C) starting temperature of the lithium capillary-porous structure (CPS), the plasma parameters were much better in comparison to the discharges with a graphite limiter. Using this limiter, it was possible to achieve a quasi-stationary mode of the discharge in deuterium at the plasma current $I_p \cong 90$ kA, shot duration ~ 0.2 s with the stabilization of the main parameters at some certain level ($\langle n_e \rangle \cong 2 \cdot 10^{19} \text{ m}^{-3}$, $T_e \approx$

400 eV, $\tau_E \approx 8$ ms, $T_{lim} \cong 260^\circ\text{C}$, $P_{lim} \approx 10$

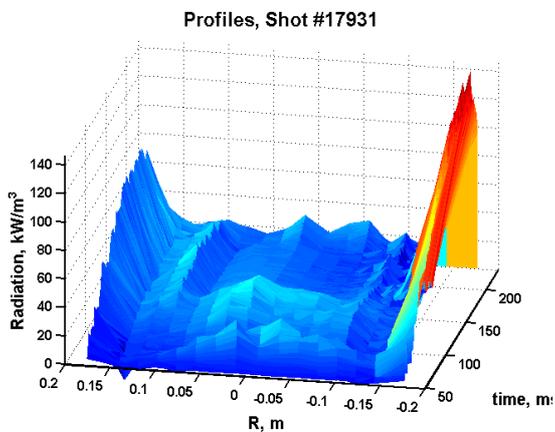


Fig. 2. The evolution of the plasma radiative loss profile with a Li limiter.

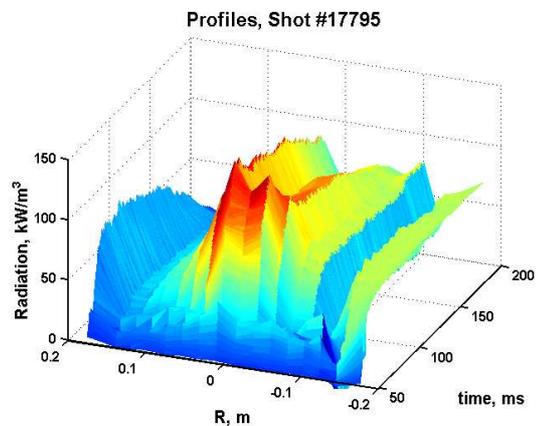


Fig. 3. The evolution of the plasma radiative loss profile with a graphite limiter.

Plasma radiative loss profile was measured by 16-channel AXUV detector array [3], installed vertically in the tangential port of T-11M tokamak [4] thus providing the toroidal

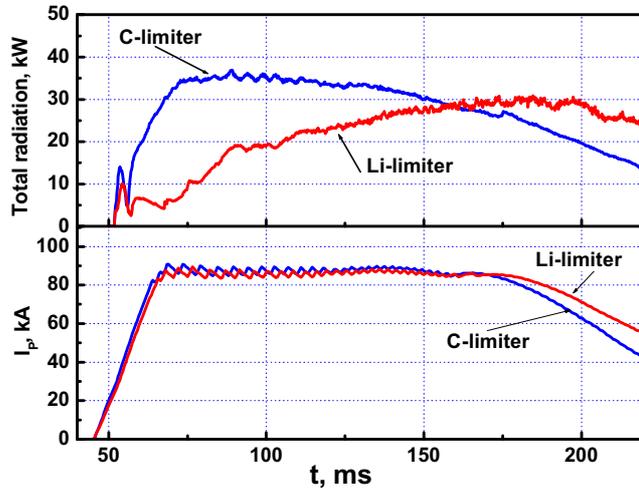


Fig.4. Total plasma radiative losses in two different discharges with Li and graphite limiters with the identical average plasma density $\langle n_e \rangle \sim 2.2 \cdot 10^{19} \text{ m}^{-3}$.

view direction. The device is able to monitor the variations of plasma emission profile in a wide spectral range with the total frame time $200 \mu\text{sec}$ and the spatial resolution $\sim 2 \text{ cm}$. The recovery of radiative loss profile was performed by the numerical procedure similar to Abel inversion adapted for the toroidal geometry [4]. The results are presented in a Fig.2 for the discharge with the lithium limiter and in the Fig.3 - for the discharge with the graphite one. The line averaged plasma density in both discharges was approximately equal $\langle n_e \rangle \sim 2.2 \cdot 10^{19} \text{ m}^{-3}$, as well as the total power of radiative losses. The

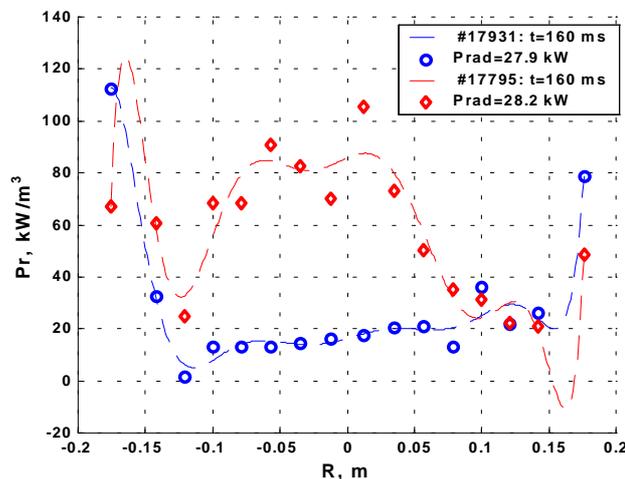


Fig.5. The radiative loss profiles at the moment $t = 160 \text{ ms}$, the average plasma density $\langle n_e \rangle \sim 2.2 \cdot 10^{19} \text{ m}^{-3}$ for both cases. Red points-C limiter, blue points -Li limiter.

comparative analysis of radiative loss profiles in these two discharges reveals the considerable difference (Fig.2, Fig.3). In the case of a graphite limiter the radiative losses are concentrated at the plasma centre (Fig.3), and in the case of a lithium one the most part of radiative losses (more than 70%) are located at the plasma edge (Fig.2). It should be emphasized, that the total plasma radiative losses and densities in both cases are approximately equal (Fig. 4). In spite of this fact, the intensity of the most undesirable radiation from the plasma centre in lithium discharge is 3-5 times (Fig. 5) suppressed in comparison to the discharges with the graphite limiter. The radiative cooling

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of the centre decreases of Ohmic heating efficiency. The $\sim 20\%$ drop of Ohmic heating efficiency of the plasma centre for graphite discharges was approximately twice more than that for lithium ones ($\sim 10\%$). The electron temperature for the lithium discharge was $T_e(0) \approx 400$ eV (Fig.1), and a bit lower - for the graphite one $T_e(0) \approx 350$ eV at a somewhat higher Ohmic power. The lower- Z_{eff} plasma in Li-limiter discharges seems to be the most probable reason of this effect. The lithium concentration at plasma centre estimated from the measured parameter $Z_{\text{eff}}/q(0) \approx 1.1$ for shots with relatively low limiter starting temperature ($\sim 250^\circ\text{C}$) is about 5%. The absolute value of radiative losses on the plasma axis was ~ 20 kW/m³ in the lithium shots, rising up to ~ 120 kW/m³ on the plasma edge. The results obtained experimentally confirm the earlier guess [2] that in the discharges with lithium limiter, the hot plasma core is enclosed within a strong radiating Li ion shell (see Fig.2). It should be noted also, that the growth of radiative losses at the plasma edge is quite helpful in the quasi-stationary tokamak discharges, providing the redistribution of a thermal loading from a limiter to the whole vacuum vessel.

Plasma edge instability at high lithium influx

At relatively low starting temperatures of lithium limiter ($250\dots 400^\circ\text{C}$) the radiation of a plasma shell is monotonically increased according to the growth of lithium surface temperature of and related increase of neutral lithium influx into the plasma (Fig.2). The

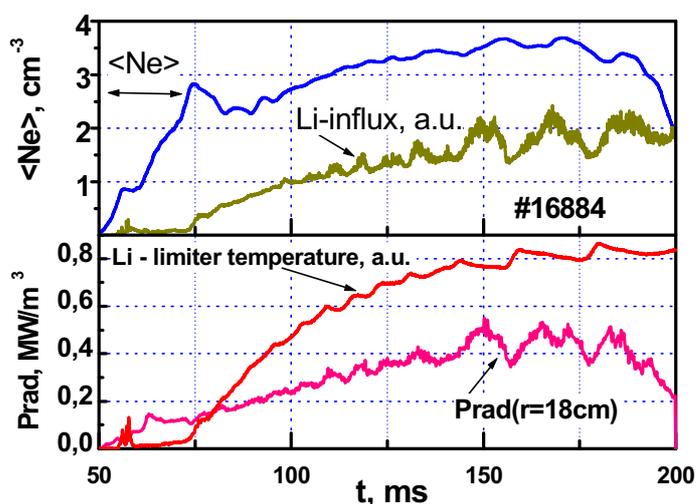


Fig. 6. Oscillations of the Li limiter surface temperature and other parameters caused by the instability at the plasma edge in the discharge with a Li limiter at high temperature $\sim 600^\circ\text{C}$.

total power of radiative losses grows similarly, since the outer plasma layers give the major contribution ($\sim 70\%$) to the total losses. However, the behaviour qualitatively varies after the transition into the higher temperature range $500\dots 600^\circ\text{C}$, and oscillations of all plasma parameters observed at the plasma edge (Fig.6,7.). Their amplitude increase following the limiter surface temperature, and the latter itself begin to oscillate

as well. The period of these oscillations are ~ 20 ms and rise and fall times, $2\dots 3$ ms and $15-$

18 ms respectively at the lithium temperature $\sim 700^\circ\text{C}$. The period of these oscillations is more at higher surface temperature. Although the nature of these oscillations is not

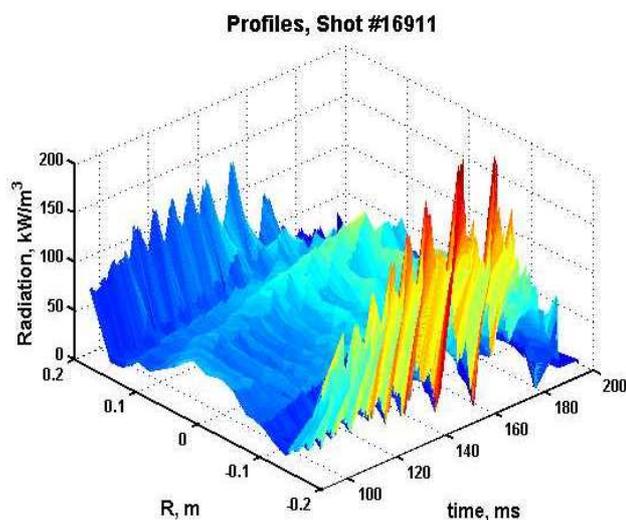


Fig. 7. The evolution of plasma radiative loss profile in the discharge with a "hot" Li limiter, temperature $\sim 580^\circ\text{C}$. Development of "ionisation-condensation" instability.

completely clear, it is possible to assume, that they are caused by the high influx of neutral lithium into the peripheral plasma area resulting to its cooling. The estimated total influx of lithium atoms from the limiter is $\Gamma \sim 5 \cdot 10^{20} \text{s}^{-1}$. Normalizing this value by the area of the outer magnetic surface of a tokamak T-11M, the average influx density is about $\Gamma/S \sim 2 \cdot 10^{20} \text{s}^{-1}/\text{m}^2$. Very close value for a threshold of transition in a non-stationary mode, or threshold of

"ionisation-condensation" instability with formation of MARFE-like region was calculated in [5], where the behaviour of lithium in SOL of the ITER-like "lithium tokamak" with the help of the two-dimensional numerical code was simulated. Both effects observed on the T-11M tokamak with lithium limiter. Namely:

- formation of a screening radiative shell at Li influx level below $\Gamma/S \sim 2 \cdot 10^{20} \text{s}^{-1}/\text{m}^2$;
- development of "ionisation-condensation" instability at Li influx above this level.

These results are in a some agreement with the behaviour of lithium in the SOL region, predicted in reference [5]. An other significant result is that such high Li influx $\Gamma \sim 5 \cdot 10^{20} \text{s}^{-1}$ do not give any disruption and the "ionisation-condensation" instability gives only oscillations on the plasma edge.

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