

Simultaneous real time control of O and F radicals in an argon/oxygen/fluorocarbon plasma

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1. Plasma process control simulation

Figure 1. is a schematic of closed loop control of a plasma process simulation. The Ar/O₂/C₄F₈ global model is adapted from the chemistry sets contained in [1], [4] and [7]. Each mass flow controller (MFC) is modelled as a first order linear transfer function, which has been determined from the step response of a commercial MFC. The relationship between the forward power P_{for} and the power P deposited in the plasma is modelled using an impedance model and a matching circuit. Transport delays, which have a destabilising effect on control loops, should be included in any simulation. Given (simulated) measurements from the plasma, the controller is required to adjust the flow rate set points F_{O2}^* and F_{C4F8}^* in order to maintain n_O and n_F at the set points n_O^* and n_F^* , respectively, despite the effects of process disturbances.

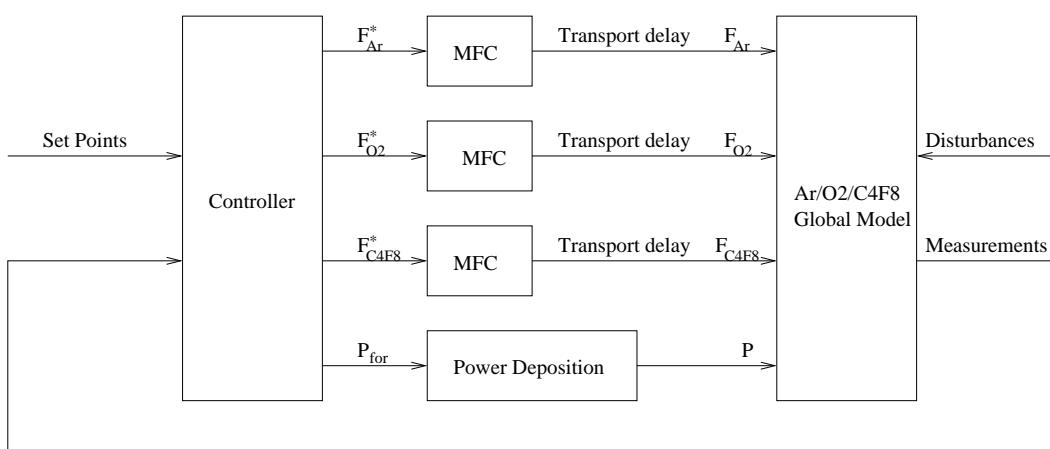


Figure 1. Plasma process control simulation

2. The transfer function for atomic oxygen density

The relationship between F_{O2} and n_O may be approximated by the linear transfer function

$$\frac{n_O(s)}{F_{O2}(s)} \approx \frac{K_O}{s\tau_{res} + 1}, \quad (1)$$

where τ_{res} is the residence time and K_O is the process gain. The relationship between F_{O2} and F_{O2}^* , may be expressed in the form

$$\frac{F_{O2}(s)}{F_{O2}^*(s)} = \frac{e^{-s\theta_m}}{s\tau_m + 1}, \quad (2)$$

where τ_m is the characteristic rise time of the MFC and θ_m is a transport delay. Equations (1) and (2) yield the process model

$$G_O(s) = \frac{n_O(s)}{F_{O2}^*(s)} = \frac{K_O e^{-s\theta_m}}{(s\tau_m + 1)(s\tau_{res} + 1)}. \quad (3)$$

3. Closed loop control of atomic oxygen

G_O defines a second order plus dead time process, which may be controlled very effectively with a Proportional-Integral-Derivative (PID) controller, the transfer function of which may be expressed in the form

$$G_c(s) = K_c \left(1 + \frac{1}{\tau_I s} + \tau_D s \right), \quad (4)$$

where K_c is the controller gain, and τ_I and τ_D are the integration and differentiation time constants, respectively. The Direct Synthesis PID settings are [5]

$$K_c = \frac{1}{K_O} \frac{\tau_m + \tau_{res}}{\tau_c + \theta_m}, \quad (5)$$

and

$$\tau_I = \tau_m + \tau_{res}, \quad \tau_D = \frac{\tau_m \tau_{res}}{\tau_m + \tau_{res}}, \quad (6)$$

where τ_c is the rise time of the desired closed loop transfer function. With τ_I and τ_D given by (6), it may be shown using the Nyquist stability criterion that the closed loop transfer function is stable if and only if $0 \leq K_c < K_u$, where

$$K_u = \frac{1}{K_O} \frac{\pi(\tau_{res} + \tau_m)}{2\theta_m}. \quad (7)$$

The gain margin, $GM = K_u/K_c$ is a measure of closed loop stability robustness, and with K_c given by (5),

$$GM = \frac{\pi}{2} (1 + \tau_c/\theta_m). \quad (8)$$

For a given θ_m , the closed loop time constant τ_c must be sufficiently large to guarantee robustness. $\tau_c = \theta_m$ yields $GM = \pi$ in (8), which is within range $GM = 1.7 - 4$ recommended in [5].

4. Estimation of atomic oxygen density from spectroscopic measurements using a Kalman filter

Actinometry offers a non-intrusive method for estimating the atomic oxygen density in oxygen-containing plasmas [3]. The emission intensity of the 777nm and 844nm lines may be expressed in the form

$$I_O \propto k_{ex}n_O + k_{dex}n_{O2}, \quad (9)$$

where k_{ex} and k_{dex} are the direct and dissociative excitation rate constants, respectively. Where the dissociation fraction is small, it is clear from (9) that a significant portion of I_O may be due to dissociative excitation. An analogous situation exists in chlorine plasmas and a possible solution to this problem was proposed in [2], where an extended Kalman filter was used to estimate atomic chlorine density. Using the same approach as [2], a Kalman filter may be designed based on the state space process model

$$\begin{aligned} \frac{dn_{O2}}{dt} &= k_{O2}F_{O2} - \frac{n_{O2}}{\tau_{res}} + w_{O2} \\ \frac{dn_O}{dt} &= k_OF_{O2} - \frac{n_O}{\tau_{res}} + w_O \\ \frac{dk_{O2}}{dt} &= w_{kO2} \\ \frac{dk_O}{dt} &= w_{kO}, \\ I_{777} &= k_{ex,777}n_O + k_{dex,777}n_{O2} + v_{777} \\ I_{844} &= k_{ex,844}n_O + k_{dex,844}n_{O2} + v_{844}, \end{aligned}$$

where $(n_{O2} \ n_O \ k_{O2} \ k_O)'$, $(w_{O2} \ w_O \ w_{kO2} \ w_{kO})'$ and $(v_{777} \ v_{844})'$ are the state, process noise and measurement noise vectors, respectively. The Kalman gain may be derived from the process model and noise statistics [6]. The process gain estimate may be used to adjust the PID controller gain according to (5), in order to maintain a sufficient gain margin.

5. Estimation and control of fluorine density

The relationship between F_{C4F8}^* and n_F may be approximated by the transfer function

$$G_F(s) = \frac{n_F(s)}{F_{C4F8}^*(s)} = \frac{K_F e^{-s\theta_m}}{(s\tau_m + 1)(s\tau_{res} + 1)}, \quad (10)$$

where K_F is a gain term. The transfer functions (3) and (10) differ only in the process gain. Hence, a PID controller with the time constants given by (6) and a gain term given by

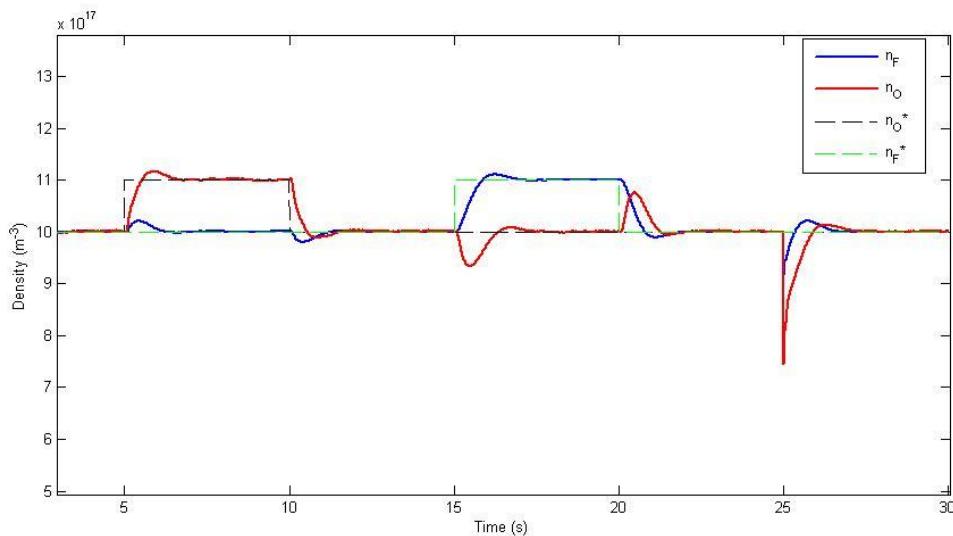
$$K_c = \frac{1}{K_F} \frac{\tau_m + \tau_{res}}{\tau_c + \theta_m}, \quad (11)$$

may be used to control n_F .

6. Results and conclusions

The closed loop responses of atomic oxygen and atomic chlorine to set point changes (shown dashed) and a wall disturbance, represented by a doubling of the atomic oxygen wall recombination coefficient

at 25s, are shown below. The model parameters used were $\tau_{res} = 50\text{ms}$, $\tau_m = 120\text{ms}$ $\theta_m = 80\text{ms}$, $\tau_c = \tau_m$. The closed loop response is seen to be stable with a slight overshoot. In practice, θ_m may be significantly larger than 80ms and systems with large time delays are difficult to control with a PID controller alone. For such systems, an approach such as dead time compensation may be appropriate [5]. In addition, it is clear that significant interaction exists between the loops. This interaction may be reduced using decoupling control [5].



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