

## Turbulent Transport Studies in Alcator C-Mod Ohmic Plasmas\*

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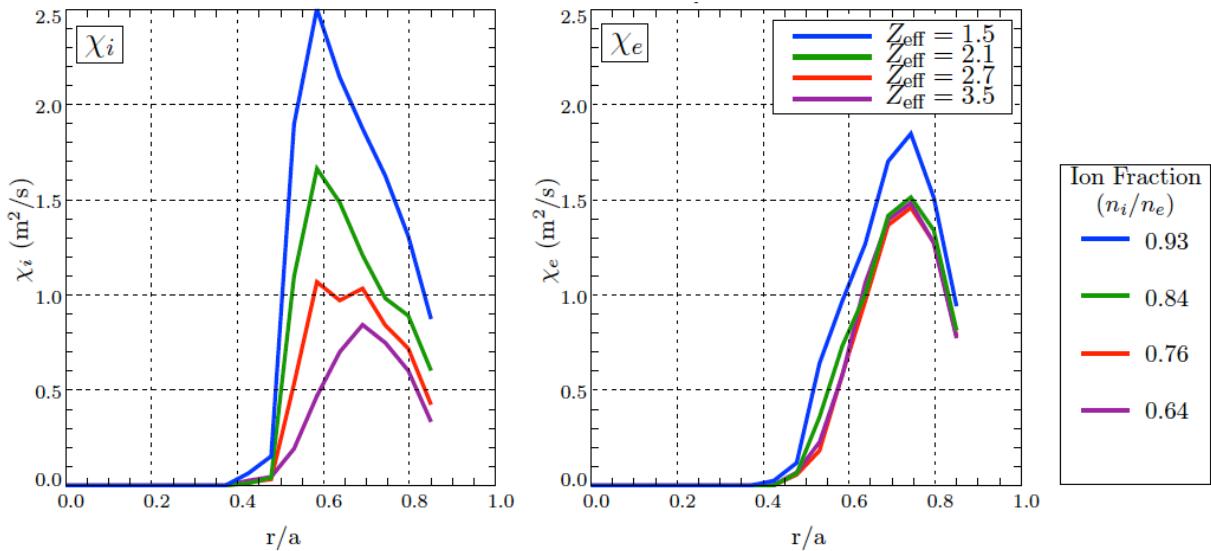
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### 1. Introduction

In our recent studies of ohmically heated low density C-Mod plasmas where  $\tau_e \propto n_e$ , the so-called neo-Alcator regime [1] where  $T_e > T_i$ , TRANSP results indicated heat diffusivities such that  $\chi_i < \chi_e$ , whereas our nonlinear gyrokinetic analysis [2] with GYRO [3] indicated that at least at low  $Z_{eff}$ ,  $\chi_i > \chi_e$  due to the dominance of ITG modes over TEM or ETG. In essence, this relatively low density regime ( $n_e \leq 1 \times 10^{20} \text{ m}^{-3}$  in Alcator C-Mod, and in fact in several other tokamak experiments worldwide) corresponds to a confinement *deterioration* relative to L-mode, in contrast to optimistic expectations in the 1980s that this confinement scaling would continue to very high densities. However, the favorable linear confinement scaling ceased at sufficiently high densities, at least in gas fuelled discharges [1], due to saturation (now we know that this occurs at the L-mode confinement value), presumably due to dominant ITG turbulence [1, 2]. In our previous experiments [2] the composition of the impurities and  $Z_{eff}$  was not well known. In addition, the ion temperature profiles were determined from TRANSP using a model  $\chi_i$  profile, calibrated to the temperature from neutron emission. These experiments were repeated recently with greatly improved diagnostics, and initial results were presented [4, 5]. In the present work the measured ion temperature profiles from x-ray crystal spectroscopy are used, and new values of the measured  $Z_{eff}$  were found to be larger than previously assumed. Recently TGLF, the trapped gyro-Landau fluid model code [6], was used with great efficiency to model the turbulent transport as a function of density and  $Z_{eff}$ , and the results indicate that the second (impurity) ion species with moderate  $Z_i$  reduces significantly the predicted values of  $\chi_i$  while maintaining the electron diffusivities near the experimental values as  $Z_{eff}$  was increased and the density decreased. By varying the impurity  $Z_i$ , it was verified that increased collisionality alone was not sufficient to explain these results. New analysis with the nonlinear flux tube GYRO [3] simulations indicates a similar trend and some of the results are presented here.

## 2. Numerical results with the trapped gyro-fluid Landau fluid model code TGLF

Here we present results with TGLF [6] which shows a significant reduction of the ion energy diffusivity as a two ion species plasma is introduced, with moderate  $Z_i$  (= 8 in the present case, corresponding to oxygen impurities) [Fig. 1]. The impurity density is varied (see dilution factors in the vertical bar) and the ensuing  $Z_{\text{eff}}$  is indicated by the different colored curves. The experimentally measured  $Z_{\text{eff}}$  corresponds to 2.7 (see red curve) at this line average density ( $n_e = 0.64 \times 10^{20} \text{ m}^{-3}$ ). We see that at the same time the dilution has negligible effect on the electron transport coefficient. Thus, we are now in a position to model the variation of transport at different densities and compare the results with the experimentally measured transport coefficients as determined by TRANSP, using the experimentally measured electron and ion temperatures, density, radiated power, and  $Z_{\text{eff}}$ . The radiated power measured in these discharges is too low for the majority of the  $Z_{\text{eff}}$  to originate from molybdenum impurities. In most of the discharges 10% or less of the  $Z_{\text{eff}}$  being from molybdenum is sufficient to produce the measured radiated power.

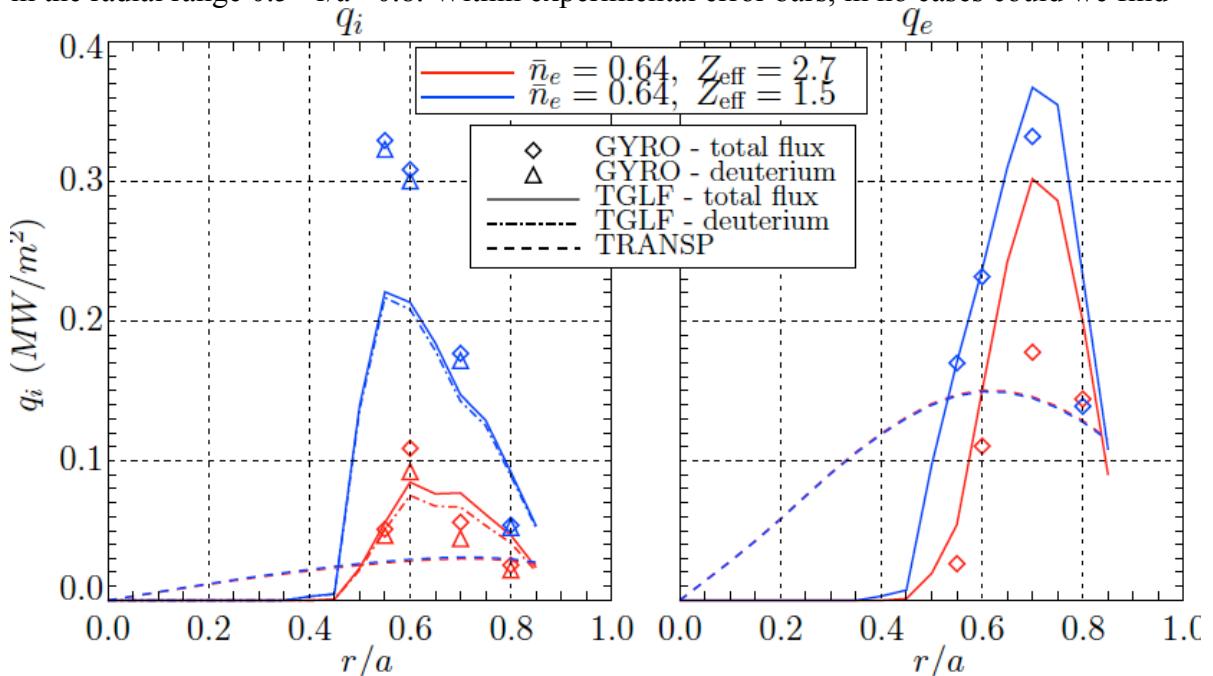


**Figure 1.** Variation of the ion and electron transport coefficients as a function of ion impurity concentration, assuming  $Z_i = 8$  (oxygen). Predictions based on TGLF. Alcator diverted ohmic discharges,  $n_e = 0.64 \times 10^{20} \text{ m}^{-3}$ ,  $B_T = 5.2 \text{ T}$ ,  $I_p = 800 \text{ kA}$ . The experimental value of  $Z_{\text{eff}}$  corresponds to 2.7, the red curve.

We see that for the experimentally measured density and temperature profiles [5], at the experimental value of  $Z_{\text{eff}} = 2.7$  (red curve) the turbulent heat diffusivities are localized to regions of maximum temperature and density gradients  $r/a = 0.5 - 0.9$ , with a peak value of  $\chi_i = 1.0$  and  $\chi_e = 1.5$ . This is a typical result, and we almost never see any significant turbulent transport in the core of the plasma  $r/a < 0.5$ , where typical experimental heat diffusivities are of the order of  $0.5 \text{ m}^2/\text{sec}$ . Note that at this  $Z_{\text{eff}}$ ,  $\chi_i < \chi_e$  as desired.

### 3. Numerical results with TGLF and GYRO

We have carried out a series of runs on the MIT PSFC parallel computing network LOKI (600 processors) with the nonlinear flux tube GYRO code at 4 different densities, including studies of the total heat flux  $q$  and heat diffusivities  $\chi$  as a function of radius, at both low and high  $Z_{\text{eff}}$  and compared the results with TGLF. Values of  $k_{\theta}\rho_s$  in GYRO are in the range 0.1-1.2, with little contribution to electron transport from higher values when tested. A plot of the radial profile of the total heat flux  $q$  is shown for both ions and electrons in Fig. 2 at a density of  $n_e=0.64 \times 10^{20} \text{ m}^{-3}$ . The TRANSP runs (experiment) are shown by the dashed lines, the blue curves and symbols correspond to the low  $Z_{\text{eff}} = 1.5$  used in the past modeling of Lin et al [2] and the red color corresponds to  $Z_{\text{eff}} = 2.7$ , which is the presently believed experimental value and used in the most recent GYRO and TGLF modeling. While the predicted values of total  $q_i$  (D and O combined) are higher than the experimental values by factors of two, using slightly modified gradients and lower  $Z_i$  (6) and by using TGYRO to iterate density and temperature profiles, reasonably good agreement between experiment and modeling is found in the radial range  $0.5 < r/a < 0.8$ . Within experimental error bars, in no cases could we find



**Figure 2.** Radial plots of the ion and electron heat flux as obtained from experiment (TRANSP, dashed curves) and TGLF (solid curves) and GYRO, (symbols) for both high  $Z_{\text{eff}} = 2.7$  (present experimental value) and  $Z_{\text{eff}} = 1.5$  (low values used in Ref 2). In the codes we assumed  $Z_i = 8$ ,  $n_e = 0.64 \times 10^{20} \text{ m}^{-3}$ ,  $B_T = 5.2 \text{ T}$ ,  $I_p = 800 \text{ kA}$ .

agreement for  $r/a < 0.5$ . Similar comparisons have been made between TRANSP predictions and TGLF and GYRO at 4 different densities in the range  $n_e = (0.35-1.05) \times 10^{20} \text{ m}^{-3}$  and the

results will be presented in a later publication. Furthermore, we have now obtained even better agreement with experiments with both nonlinear flux tube GYRO, and TGLF iterated with TGYRO, a technique which allows the profiles to be adjusted to match to  $q_i$  and  $q_e$ , the ion and electron power flux. One question in such modeling is the role of stabilization of the ITG modes by shear flow. Recent measurements in similar discharges [7] indicated a maximum  $E_r$  of the order of 14 kV/m. Using the measured  $E_r$  radial profiles, we found at most a 10% reduction in the predicted ion transport.

### 3. Summary and Conclusions

Extensive gyrokinetic studies of low density ohmic plasmas in Alcator C-Mod revealed the important role of deuterium depletion in the presence of a second (impurity) ion species with moderate  $Z_i$ , resulting in reduced ion transport and dominant electron transport in the outer half of the plasma column. Additional examination of linear growth rates is in agreement with this observation. The radial electric field shear does not play a dominant role in the stability of these plasmas. Transport in the inner half of the plasma core is not explained by the gyrokinetic codes. Potential physical mechanisms to give the enhanced transport in the plasma core include mild sawteeth (although the inversion radius is typically  $r/a \leq 0.35$ ), or drift waves driven by ohmic electron drift which at the lowest densities considered here can attain values as high as  $U/C_s = 6$  near the center, and diminish radially in accordance with the electron temperature profile [4].

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