

## Nonlinear saturation of the Toroidal Alfvén Eigenmodes computed with the VENUS+δf, HAGIS and KINX codes

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### 1. Introduction

The successful benchmarking of the linear growth stage of fast particle driven TAE with a broad variation of numerical models by the ITPA Energetic Particle Topical Group for an  $n = 6$  Toroidal Alfvén Eigenmode (TAE) number was described in [1]. In this paper, the second stage in the wave evolution – the non-linear saturated state – is computed for the same  $n = 6$  test case (Section 2) and also for an  $n = 4$  TAE in JET #40214 discharge (Section 3). The fixed spatial structure of TAE modes  $\eta_{mn}(s)$  ( $s = (r/a)^2$  is in this paper a normalised toroidal (or can be also a poloidal in some codes) flux function,  $a$  is a minor plasma radius) is calculated with the KINX [2] and MISHKA codes [3]. The fast particle dynamics, TAE growth rates and wave saturation levels are computed with the HAGIS [4] and VENUS+δf [5, 6] orbit following codes.

Fourier decomposition in Boozer coordinates  $(s, \theta, \zeta)$  of the TAE mode with a poloidal index  $m$  and the mode frequency  $\omega$  has a form

$$\xi = A(t) \sum \eta_{mn}(s) \cos(m\theta - n\zeta - \omega t). \quad (1)$$

The main equation for the TAE amplitude  $A(t)$  evolution due to the fast particle-wave interaction without plasma damping is

$$dA/dt = - \langle \int Z e \delta f \mathbf{V} \cdot \mathbf{E} d\tau \rangle / (2K \omega^2 A), \quad (2)$$

here  $K = \int \rho_b \xi^2 dV$  – the kinetic energy of the plasma perturbation,  $\rho_b = m_b n_b$  – mass density of the bulk plasma, volume unit,  $d\tau = d^3x d^3V$  – phase volume unit,  $Z \cdot e$  – particle charge,  $\delta f$  – perturbed distribution function of fast particles,  $\mathbf{V}$  – particle velocity vector,  $\mathbf{E}$  – wave electric field. The TAE growth rate  $\gamma$  is computed from  $\gamma = dA/(A dt)$ , and the radial component of the TAE perturbation  $\delta \mathbf{B} = \text{rot} \alpha \mathbf{B}$ , normalized to the central magnetic field  $B_0$ , is defined as

$$\delta B_r/B_0 = \delta \mathbf{B} \nabla s / [|\nabla s| B_0] = -\mu_0 (I \partial \alpha / \partial \theta + J \partial \alpha / \partial \zeta) / [|\nabla s| g^{1/2} B_0], \quad (3)$$

where  $J$  and  $I$  are the toroidal and poloidal current flux functions,  $g^{1/2}$  is a Boozer jacobian.

## 2. Simulation results for the ITPA-EP test case with $n = 6$ TAE

A circular tokamak with major and minor radii  $R = 10$  m,  $a = 1$  m respectively, safety factor  $q(s) = 1.71 + 0.16 s$ , and magnetic field on-axis of 3T was chosen as an ITPA-EP test case because of the restrictions of the 9 participating codes with respect to geometry or numerical properties. A detailed description of this test case is provided in Ref. [1]. The plasma cross-section with the  $n = 6$  TAE perturbation computed with the KINX code, is shown in Fig. 1. The successful benchmark of the linear TAE growth with a Maxwellian distribution of fast ions (deuterons) with a temperature range of  $T = 100 - 800$  keV has been performed with both zero and finite Larmor radius effects. The maximum linear growth rate was  $\gamma_L = 5 \times 10^4 \text{ s}^{-1}$ . The typical nonlinear saturation of the TAE mode function  $\delta B_r/B$  for  $T = 400$  keV, computed with the VENUS code, is shown on Fig. 2 in green. The evolution of the growth rate is shown in blue. Good agreement between the VENUS nonlinear simulations and a theory scaling  $\delta B_r/B \sim (\gamma_L)^2$  [7] is shown in Fig. 3.

## 3. Simulation results for the JET#40214 discharge with TAE $n = 4$

JET DD discharge #40214 at  $t = 46.38$  s has been selected for making more realistic simulations as it has an elongated plasma cross-section, moderate aspect ratio and a core-localised  $n = 4$  TAE mode with a frequency  $\omega = 1161100 \text{ s}^{-1}$  (see Fig. 4 which shows the plasma cross-section and TAE perturbations computed with the KINX code). The bulk is taken to be a DT plasma (50% of deuterium and 50% of tritium) with a central plasma density  $n_0 = 2 \times 10^{19} \text{ m}^{-3}$  and a flat ion density profile  $n = n_0 (1 - 0.15 \Psi^2 - 0.85 \Psi^4)$ , where  $\Psi$  is the normalised poloidal flux and the ion temperature is constant,  $T = 15$  keV. The fast  $\alpha$ -particles are assumed to have an exponential density profile  $n_f = n_{f0} \exp(-5.5 \Psi)$  and a slowing down distribution function in energy,  $f(E) = C(1 - \text{erf}(x))/(E^{3/2} + (E_c)^{3/2})$  with  $E_c = 4.942 \times 10^5 \text{ eV}$ ,  $E_0 = 3.5 \times 10^6 \text{ eV}$ ,  $dE = 4.105 \times 10^5 \text{ eV}$ ,  $x = (E - E_0)/dE$ . The constant  $C$  has been defined to provide a central  $\alpha$ -particle density of  $n_{f0} = 6 \times 10^{16} \text{ m}^{-3}$ .

Fig. 5 shows the nonlinear saturation of the  $n = 4$  TAE mode computed with the VENUS code for the JET #40214 discharge with different initial amplitudes:  $A(0) = 1.e-9, 1.e-8, 1.e-7, 3.e-7, 5.e-7, 7.e-7$ . The saturation level  $\delta B_r/B \approx 2 \times 10^{-5}$  does not depend on this initial value, however, the saturation can be achieved much faster (after about 200 wave periods) with a large value of  $A(0) = 7.e-7$ . These VENUS computations use about  $10^6$  particles, time step  $5.e-8$  s. The saturation level  $\delta B_r/B \approx 2 \times 10^{-5}$  has been achieved with the HAGIS code with  $10^6$  particles after 800 wave periods (Fig. 6). The linear stage has the

normalized growth rate  $\gamma_L/\omega \approx 0.15\%$  from both the HAGIS and the VENUS computations, however, for this core-localized TAE case VENUS results depend on the complicated particle orbits near the magnetic axis.

#### 4. Summary and future plans

After the successful benchmark for the linear stage of the TAE evolution in the frame of ITPA-EP group, we present the simulation results of the nonlinear evolution of the TAE modes with a fixed spatial structure and phase. Nonlinear TAE saturation levels for the ITPA-EP  $n = 6$  test case computed with the VENUS+ $\delta f$  code are in an agreement with the theory scaling.

Nonlinear saturation levels for JET #40214 discharge with an  $n = 4$  TAE mode, computed with the HAGIS and VENUS+ $\delta f$  codes, are in a good agreement and equal to  $\delta B_r/B \approx 2 \times 10^{-5}$ . Further benchmarks will be performed in the frame of ITPA-EP group. VENUS code will explore the wide orbit effects near the plasma edge and near the magnetic axis, where the VMEC code can have a poor equilibrium force balance [8]. Extended JET-DT TAE experimental details, phase evolution equation, damping effects and relaxation effects (sink, source, diffusion or drag) will be considered in the future to explore predictions for TAE behavior in ITER.

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#### References

- [1] A. Könies *et al*, 24<sup>th</sup> IAEA Fusion Energy Conference, CN-197, ITR/P1-34(2012).
- [2] L.Degtyarev *et al*, Comp. Phys. Comm., 103, 10(1997).
- [3] A.B.Mikhailovskii *et al*, Plasma Phys. Rep., 23, 844(1997).
- [4] S.D.Pinches *et al*, Comp. Phys.Comm., 111, 133(1998).
- [5] W.A.Cooper *et al*, Plasma Phys. Contr. Fus. 53, 024001(2011).
- [6] M.Yu.Isaev *et al*, Plasma and Fus. Res., 7, 1403077(2012).
- [7] H.L.Berk, B.N. Breizman, Phys. Fluids B 2, 2226(1990).
- [8] R.Sanchez *et al*, Comp. Phys.Comm., 141, 55(2001).

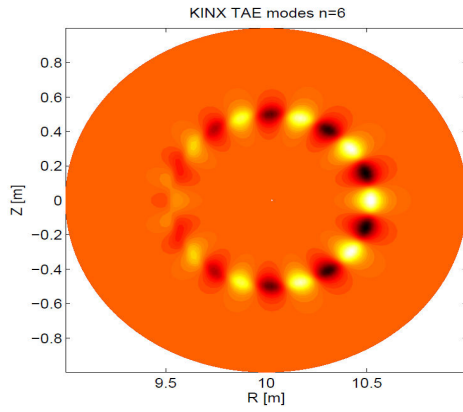


Fig. 1. Plasma cross-section with  $n = 6$  TAE perturbation computed with the KINX code for the ITPA-EP test case.

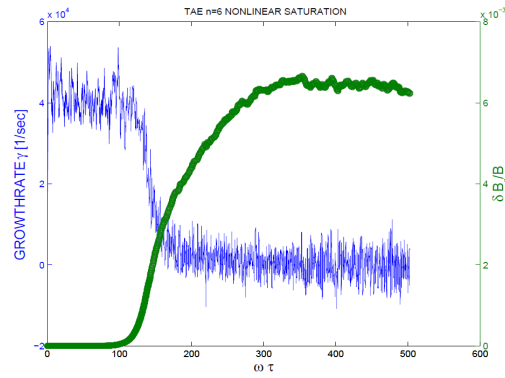


Fig. 2. Nonlinear saturation of  $n = 6$  TAE mode (green), computed with the VENUS code for the ITPA-EP test case. Evolution of the TAE growth rate is shown in blue.

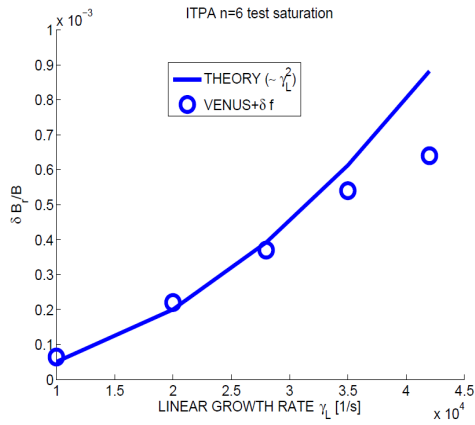


Fig. 3. Saturation level  $\delta B_r/B$  of TAE  $n = 6$  mode as a function of the linear growth rate  $\gamma_L$ , computed with the VENUS code and from a theory scaling [7].

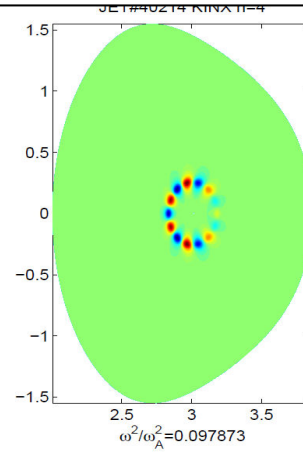


Fig. 4. Plasma cross-section with  $n = 4$  TAE perturbation computed with the KINX code for the JET #40214 discharge.

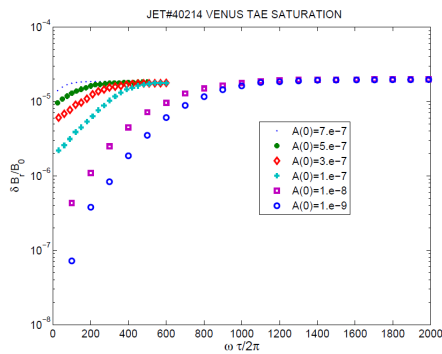


Fig. 5. Nonlinear saturation of  $n = 4$  TAE mode, computed with the VENUS code for the JET #40214 discharge with the different initial amplitudes,  $A(0)$ .

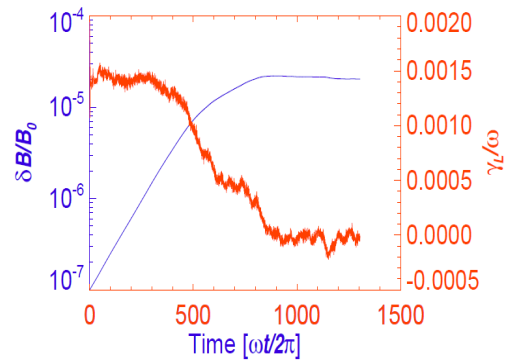


Fig. 6. Nonlinear saturation of  $n = 4$  TAE mode (blue), computed with the HAGIS code for the JET #40214 discharge. Evolution of the TAE growth rate is shown in red.