

Discharge system consisted of electrodes, voltage source, storage capacitance $C_1 = 0,22 \mu\text{F}$ ($C_2 = 0,011 \mu\text{F}$), and low inductance currents conductor. To increase the intensity of x-rays emission the conical configuration was applied [4]. Conical cathode was chosen as a high-voltage electrode, and anode with a tip diameter of $250 \mu\text{m}$ was a target-electrode where laser beam was focused. Electrodes were arranged at an angle ($\alpha \sim 30^\circ$) from each other to create the opportunity for simultaneous study of X-rays emission and ion flux (fig.2).

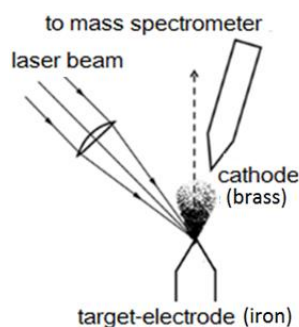


Fig.2. Principal experimental scheme.

In the first series of experiments the discharge current was sustained with a capacitor of $C_1 = 0,22 \mu\text{F}$ loaded up to -13 kV . The level of energy input into the discharge reached $E_I \sim 17 \text{ J}$ and the magnitude of discharge current didn't exceed a value $I_I = 10 \text{ kA}$. The distance between electrodes was chosen as varied parameter, which affected spectral composition and magnitude of x-rays, created in plasma jet. The size of the electrode gap was changed from 2 mm to 7 mm with step $\sim 1 \text{ mm}$.

Analysis of the spectral composition of the x-rays was performed using a series of debilitating Al filters with the range of thickness $h = 2 \div 120 \mu\text{m}$. The set of Al filters and X-rays film was placed at a distance $l_f = 6 \text{ cm}$ from the radiation source inside protective shielding.

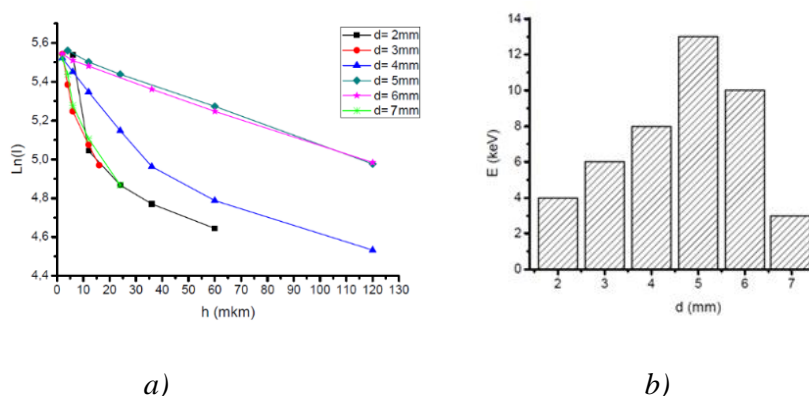


Fig.3. a) X-ray attenuation curves for filters, b) Spectral composition of the X-ray source according to the distance between the electrodes.

Figure 3(a) demonstrated that the widest range of energies and the highest intensity of X-rays emission was observed at a size of electrode gap $d = 5 \div 6 \text{ mm}$. Spectrum of X-rays was calculated by mathematical method of “effective energy” [5], (fig.3(b)).

In the second series the discharge current was sustained with a capacitor of $C_2 = 0,011 \mu\text{F}$ (fig.4). The level of energy input into the discharge was decreased to the value of $E_2 \sim 1 \text{ J}$ and the magnitude of discharge current reached $I_2 = 2 \text{ kA}$. Other parameters were not changed. The most intense radiation was observed at an electrode gap of $d = 2\div 3 \text{ mm}$.

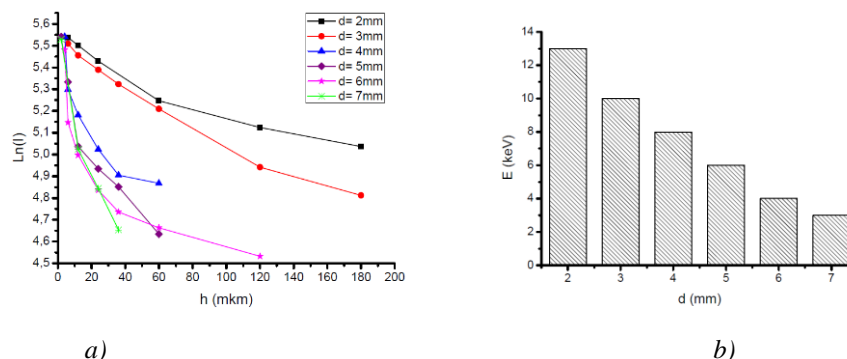


Fig.4. a) X-ray attenuation curves for filters, b) Spectral composition of the X-ray source according to the distance between the electrodes.

To obtain the spatial distribution of x-ray emission the automatic vacuum pinhole camera was applied. It located inside the vacuum vessel at the distant of 6 cm perpendicular to discharge axis. The aperture size was $d_{ap} = 500 \mu\text{m}$. It was covered by thin Al foils $h = 4\div 12 \mu\text{m}$ thickness. The images of the discharge gap were detected on the x-ray film *Kodak* (fig.5). It should be noted that bright images of plasma was obtained only for the thin Al filter $h = 4 \text{ mm}$.

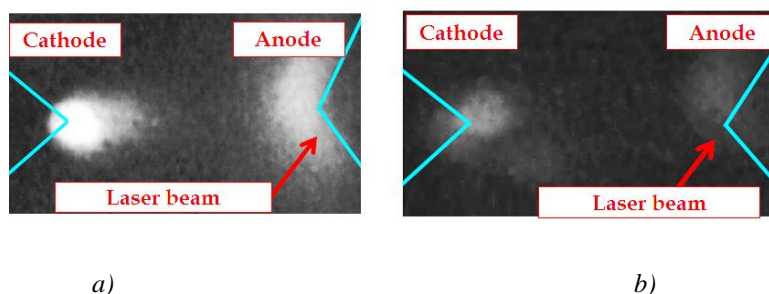
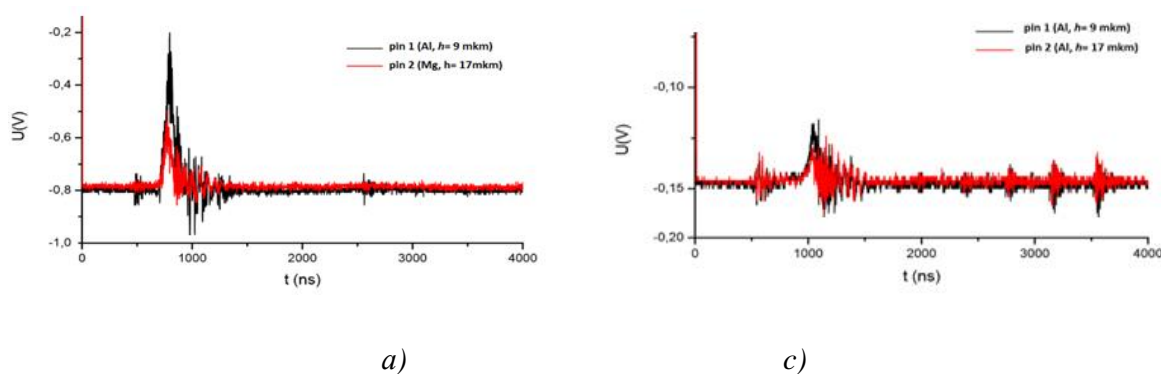


Fig.5. The pinhole images of the interelectrode gap, size $d = 5 \text{ mm}$, $C_1 = 0,22 \mu\text{F}$, $U = -13 \text{ kV}$
 a) Al filter $h = 4 \mu\text{m}$, b) Al filter $h = 12 \mu\text{m}$

To control the temporal dependence of discharge current the Rogowsky coil was placed in the cathode circuit. The temporal characteristics of X-rays emission were measured by a system consisting of two pin-photodiodes that were covered by filters (fig.6).



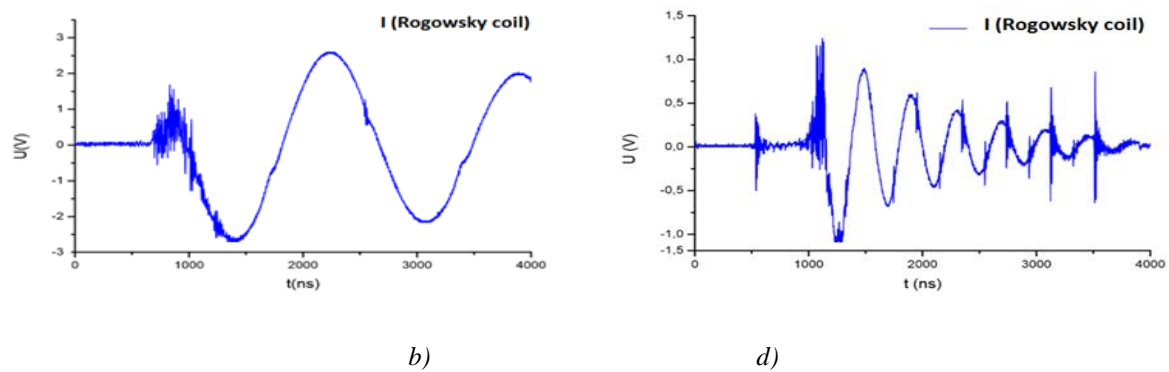


Fig.6. The dynamics of the laser induced vacuum discharge

a) Pin-photodiodes signal, $C_1 = 0,22 \mu\text{F}$, $E_1 \sim 17\text{J}$; b) Oscillogram of discharge current, $C_1 = 0,22 \mu\text{F}$, $E_1 \sim 17\text{J}$; c) Pin-photodiodes signal, $C_2 = 0,011 \mu\text{F}$, $E_2 \sim 1\text{J}$; d) Oscillogram of discharge current, $C_2 = 0,011 \mu\text{F}$, $E_2 \sim 1\text{J}$.

It is shown that moment of emergency the discharge current oscillations correlated with a time when x-ray radiation has been emitted. The discharge current has damped oscillatory in the process of developing of plasma. It's necessary to note that X-ray radiation was emitted in the initial stage of the discharge firing after $t \sim 300$ ns following laser initiation pulse. Duration of X-rays pulse was estimate $t_{xr} \sim 100$ ns and total energy of X-rays in pulse reached a value $D_1 \approx 0,4$ mJ (for $C_1 = 0,22 \mu\text{F}$, $E_1 \sim 17\text{J}$) and $D_2 \approx 2$ μJ (for $C_2 = 0,011 \mu\text{F}$, $E_2 \sim 1\text{J}$) in full solid angle.

CONCLUSION

The new kind of adjustable source of x-ray emission based on laser induced moderate power vacuum discharge with a range of quantum energy $h\nu = 1 \div 12$ keV was created. The possibility to adjust the X-ray spectral composition by changing the geometry of the electrode system was demonstrated. In the case of using stored energy $E_1 = 17\text{J}$ the highest X-rays emission was observed when $d = 5$ mm and the spectral composition consisted essentially of hard component of radiation. The intensity decreased monotonically with grows of energy from 1 to 12 keV. In the case of $E_2 = 1\text{J}$ maximal intensity was obtained when $d = 2 \div 3$ mm.

LITERATURE

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