

# Stellarator divertor optimisation for a Stable Quasi-Isodynamic Design (SQuID): magnetic topology, divertor plates and baffle design

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## Introduction

In this work we study stellarator divertor optimisation for a "Stable Quasi-Isodynamic Design" (SQuID). SQuIDs are a promising new class of optimised stellarator, receiving significant attention from the public and private sector [1,2,3]. "Quasi-Isodynamic" refers to the method of minimising neoclassical transport and "stable" refers to the fact that these stellarators are linearly MHD stable up to a sufficient plasma  $\beta$  and have relatively low electrostatic turbulent transport compared with existing stellarators. SQuIDs have also been shown to be compatible with buildable magnetic coils [2,3].

Stellarator divertors must address several challenges, for example: heat loads on plasma-facing components (PFCs) and plasma temperatures at the PFCs must be within material limits, and Helium ash and impurities must be efficiently removed to avoid core pollution. To date, the most well-explored stellarator divertor concept, and a natural candidate for SQuIDs, is the island divertor, successfully used by Wendelstein 7-X (W7-X) [4], in which an island chain diverts heat and particles away from the confined region onto divertor plates. Here we examine the edge magnetic topology of and create a first iteration of PFC design for a particular SQuID [2], a four field period device with 40 modular coils (5 distinct shapes, repeated and rotated eight times), a major radius of 20m, an average on-axis magnetic field strength of 7.5T, a core plasma  $\beta$  of  $\beta_{\text{core}}=8\%$  and 3GW modeled fusion power. A top-down view of the coils and plasma is shown in Figure 1 (left), and the magnetic structure at several toroidal locations at a single plasma  $\beta_{\text{core}}=8\%$  is shown in Figure 1 (right).

## 2. Magnetic structure

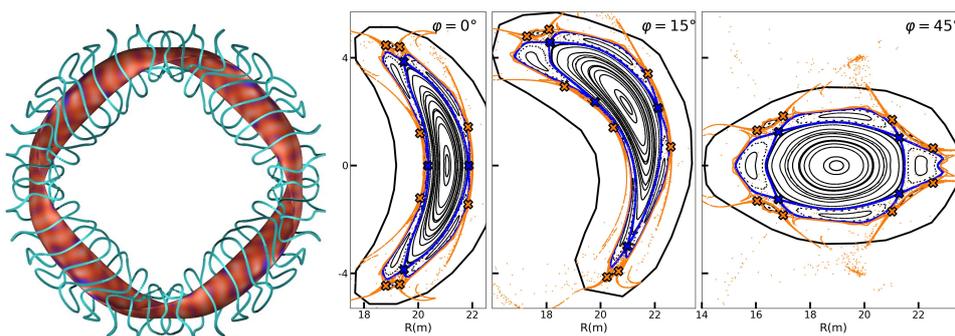


Figure 1. Top-down view of SQuID and magnetic structure (Poincaré and edge structures) at several toroidal locations at  $\beta_{\text{core}}=8\%$ .

Two edge structures are visible in Figure 1: a conventional 4/4 island chain (blue), and a second set of X-points and O-points (orange). These differ from a normal island chain because there are eight X-points but only four O-points (verified using the scheme presented in [5]). Figure 2 shows these structures at toroidal location  $\varphi=45^\circ$  at several  $\beta$  values, calculated using the HINT code [6]. The conventional island chain changes as  $\beta_{\text{core}}$  increases from 0% to 8%: from 4/4 to 8/8, and then back to 4/4 with reversed island phase (X-point  $\rightarrow$  O-point). The second (orange) structures (more precisely, the manifolds of the X-points, see [5] for explanation) appear resilient to  $\beta$ , which could be attractive for divertor design. We refer to these as “tentacles” and examine in section 3 whether they are divertor-relevant.

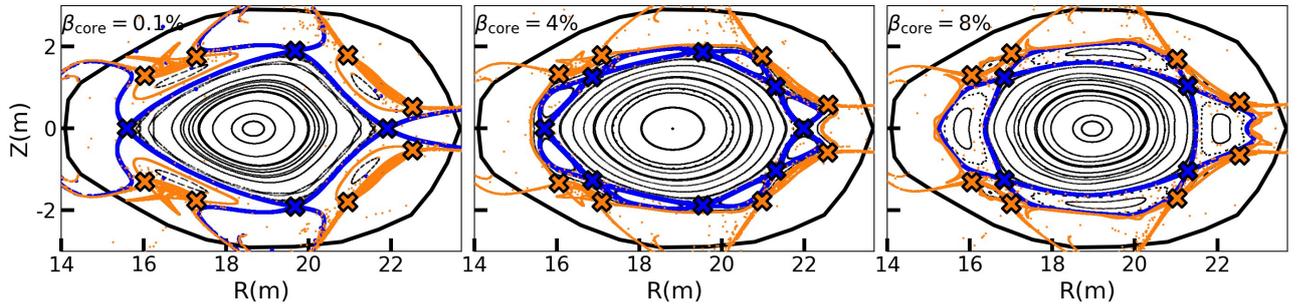


Figure 2. Magnetic structure at several values of plasma  $\beta$ .

Section 3 only considers PFC designs for  $\beta_{\text{core}}=8\%$ . However, because island phase transition is likely to change the heat and particle strike location, we demonstrate that the phase at  $\beta_{\text{core}}=0$  can be “corrected” to match that at  $\beta_{\text{core}}=8\%$ , by changing the coil currents but *without* changing the coil geometry. We do this (a Poincaré section and edge structures are shown in Figure 3), by scanning the coil currents and detecting fixed points using a fast automated scheme, which will be described fully in [7]. Changing the coil currents, will affect

### 3. PFC design

We design divertor plates using the scheme presented in [8], with simulated heat loads calculated using an anisotropic diffusion model implemented in EMC3-Lite [9]:  $\nabla \cdot (-\kappa_e \nabla_{\parallel} T - \chi n \nabla_{\perp} T) = 0$  (where  $n$ ,  $T$  are the plasma density, temperature,  $\kappa_e = \kappa_{e0} T_{LCFS}^{5/2}$  is the parallel thermal conductivity,  $\kappa_{e0}$  is a constant and  $T_{LCFS}$  is  $T$  at the last closed flux surface and  $\chi$  is the perpendicular diffusivity).

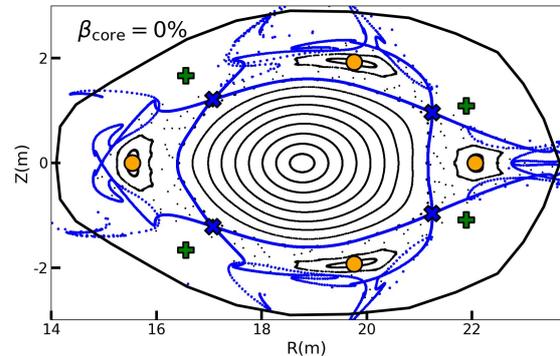


Figure 3. Corrected island phase at  $\beta=0$ .

the neoclassical transport, MHD stability and turbulent transport, although we do not investigate this here. Future work includes designing coils to ensure a  $\beta$ -resilient island phase.

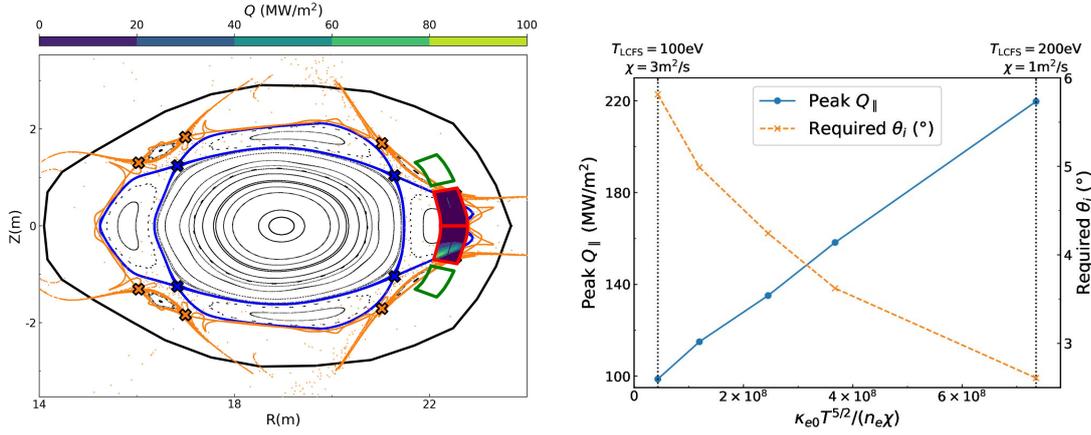


Fig. 4. PFCs for an island divertor. Left: Divortor plates (outlined red) and baffles (outlined green), before being “stretched” along the magnetic field to reduce heat loads. Heat loads shown for  $T_{LCFS}=100\text{eV}$ ,  $n_{LCFS}=10^{19}\text{m}^{-3}$ ,  $\chi=3\text{m}^2/\text{s}$ . Right: Peak  $Q_{||}$  as a function of parallel-to-perpendicular diffusivity and maximum  $\theta_i$  which ensures  $Q < 10\text{MW}/\text{m}^2$ .

Firstly, “vertical plates” are constructed at a single toroidal location, designed to catch a sufficiently high fraction of non-radiated heat. Secondly, the plates are “stretched” along the magnetic field, reducing the incidence angle  $\theta_i$  and increasing the wetted area to reduce PFC heat loads  $Q$ . Two important parameters here are: (1) the assumed non-radiated power, which we take to be  $60\text{MW}$ , consistent with  $3\text{GW}$  fusion and  $90\%$  radiated fraction; (2) the ratio of parallel-to-perpendicular diffusion,  $\kappa_e/(n_e\chi)$ , which determines the required  $\theta_i$  assuming the simple relationship  $Q = Q_{||}\sin(\theta_i)$  where  $Q_{||}$  is the parallel heat flux, and is  $\theta_i$ -independent.  $\kappa_e/(n_e\chi)$  is not known *a priori*; we therefore scan a range of values. Locations where baffles can be placed are also found, using the methodology presented in [10]; we seek locations with low total heat loads and low  $Q_{||}$ . These are intended to impede the neutral particles, promoting neutral and plasma density build-up at the PFCs.

Figure 4 shows an island divertor for the SQUID; the vertical plates (left subplot) intersect the  $4/4$  island chain, and are then stretched along the magnetic field to ensure  $Q < 10\text{MW}/\text{m}^2$  for all  $\kappa_e/(n_e\chi)$ . Baffle locations are also shown. The peak  $Q_{||}$  and required  $\theta_i$  as a function of  $\kappa_e/(n_e\chi)$  is shown in Figure 4, right. We also find vertical plates which intersect the tentacles and catch over  $90\%$  of the non-radiated heat. To do so, it is necessary to place targets at 2 toroidal locations,  $\varphi=0^\circ$  and  $\varphi=45^\circ$  (shown in Figure 5).

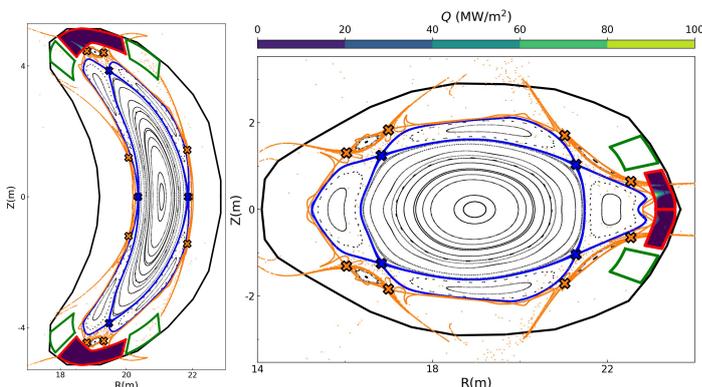


Fig. 5. Divortor plate (red) and baffles

(green) for “tentacle” divertor “vertical plates”.

## **5. Conclusion and outlook**

This work describes magnetic topology and PFC design for a SQUID. We show a geometrically “open” island divertor, with divertor plates intersecting a 4/4 island chain. The island chain is sensitive to plasma  $\beta$ ; we present a crude phase correction by changing the coil currents.  $\beta$ -resilient island design is future work. We also find  $\beta$ -resilient “tentacles”, which appear compatible with divertor plates. Our PFC design scheme, by itself, is insufficient to construct a provably reactor-relevant divertor solution, but serves as a starting point.

The simple transport model used allows fast iterations of PFC design but is the most serious drawback of our scheme. Plasma transport, radiation and neutral dynamics are likely to be complex in a stellarator reactor and require comprehensive high-fidelity simulations. A hierarchy of tools, from cheap but simple to comprehensive but expensive, is likely necessary for reactor-relevant divertor optimisation.

## **Acknowledgments**

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